

Weak and Strong turbulence in MHD

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Aim: find numerical proof of weak and strong regimes for strongly anisotropic turbulence, assumed to be valid for strong mean field B°

Motivation: when large scales are forced weakly (with $k_{\parallel}/B^{\circ} > k_{\perp}u$), one should find a weak regime (k_{\perp}^{-2}), followed by a strong regime ($k_{\perp}^{-5/3}$) at smaller scales.
No numerical proof of that exists yet.

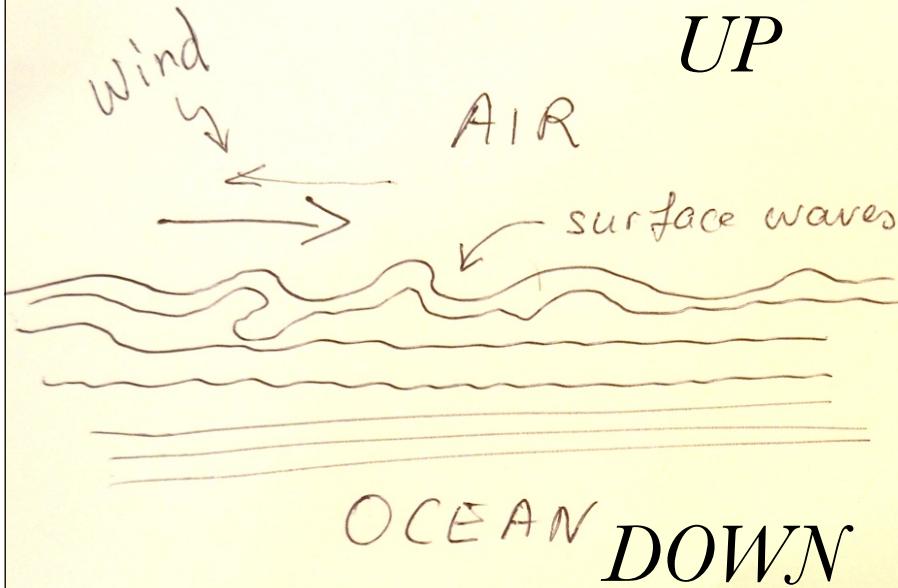
Published works report that varying the strength of large scale forcing leads to a continuous variation of the slope, but nothing like the above double scaling.
(eg Perez & Boldyrev 2008)

The shell-Reduced MHD model (Buchlin, 2004 thesis) offers the opportunity to test these ideas as it is able to describe Reduced MHD turbulence with $\text{Reynolds} \approx 10^6$

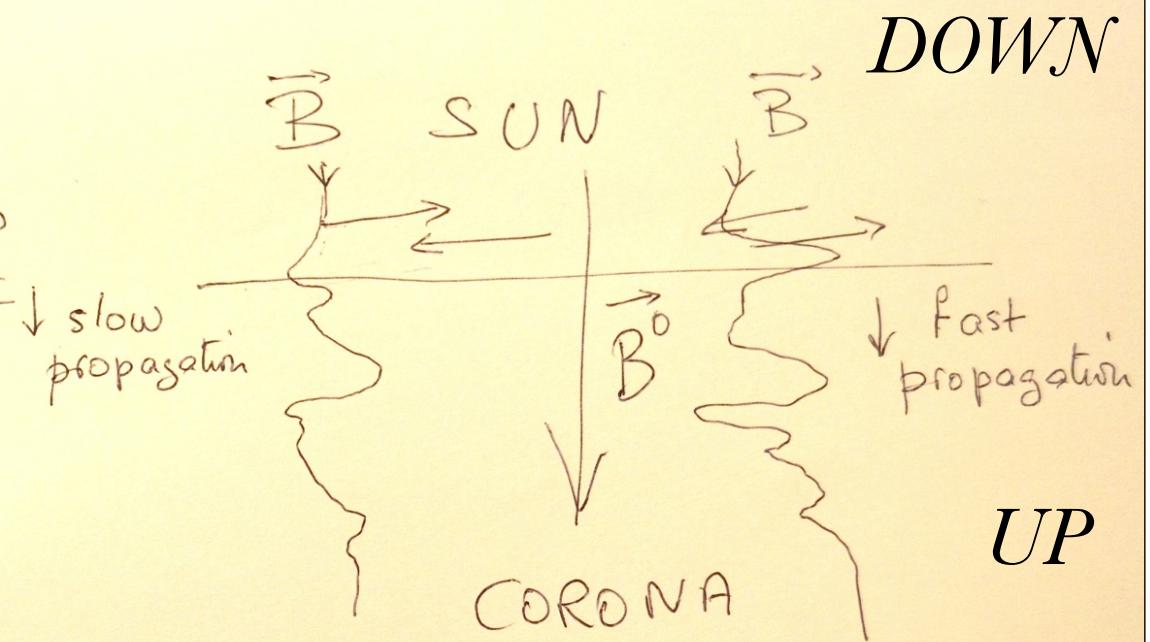
MHD with mean field B°
allows fast transmission of surface forcing



Slow transmission (no B°)



Fast transmission ($B^\circ \neq 0$)



Physical/Numerical model

1. Reduced MHD ($b/B^{\circ} \ll 1$)

- No parallel gradients except linear (propagation) along B°
 - No parallel polarizations ($b_{//}=u_{//}=0$)
 - Incompressible limit ($\nabla_{\perp} \cdot b_{\perp} = 0$)
- => **Quasi 2D** ($\perp B^{\circ}$), but ***not completely***: $u, b = u(x), b(x)$

2. "Shell" Reduced MHD (*Buchlin, thesis 2004, Nigro et al 2004, Buchlin Velli 2007*)

- solves for $\hat{u}(x, k_y, k_z)$ (1/2 FFT)
- use scalar (wavenumber) to represent \perp Fourier space by shells: $k_n \approx 2^n k^{\circ}_{\perp}$
- allows to reach $Re \approx 10^6$

Forcing strong or weak regimes

Forcing "strong" or "weak" \rightarrow cascade strong or weak

(1) "Strong" means that z^+ and z^- wavepackets have time to couple:

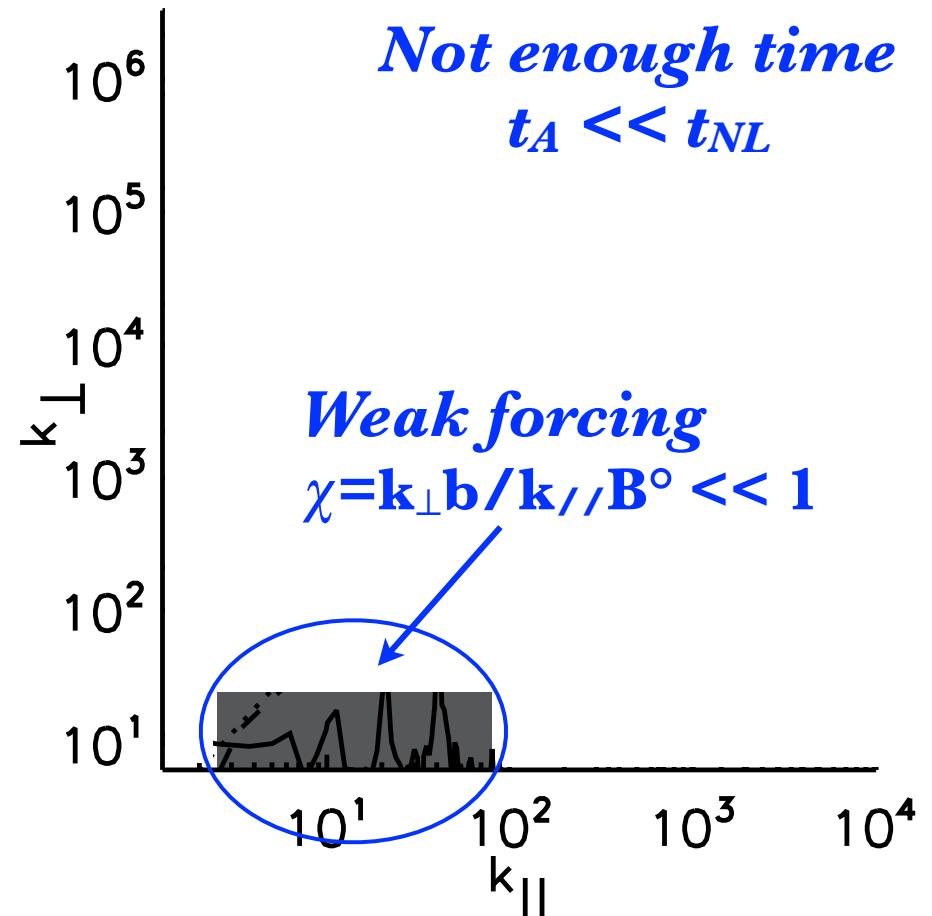
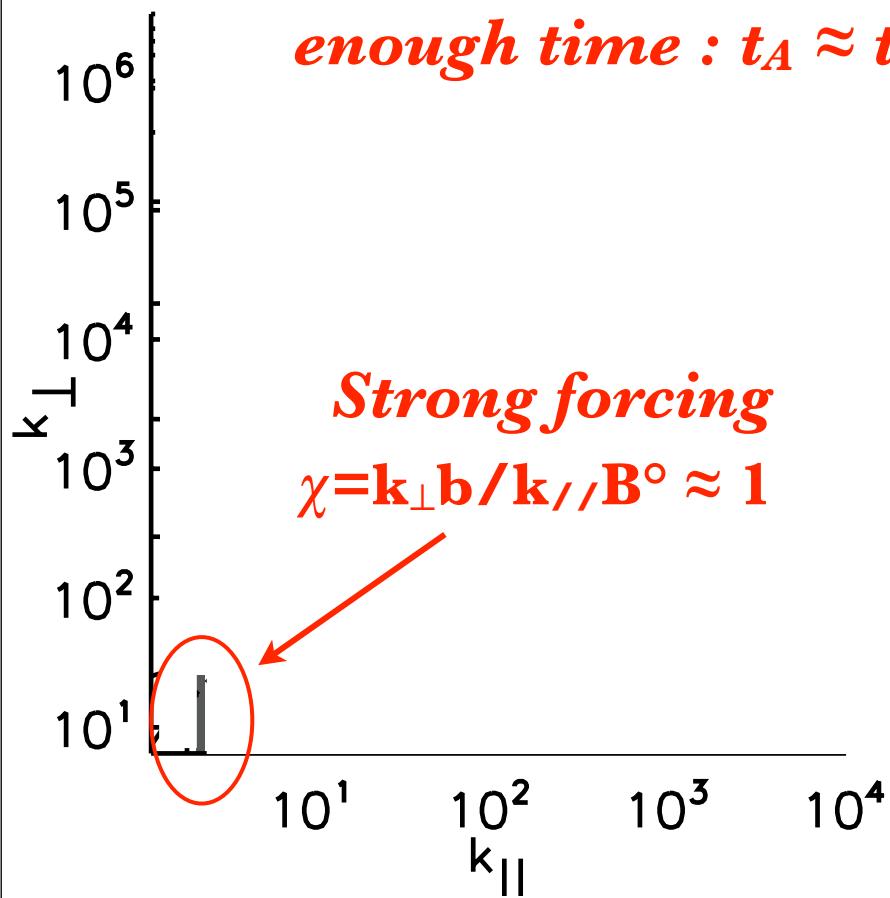
$$t_A \approx t_{NL}$$

(2) "Weak" means wavepackets pass too fast to interact completely

$$t_A \ll t_{NL}$$

Depends on ratio of $\chi = t_A / t_{NL} = kz^\pm / k_{\parallel} B^0$ at **forcing scales**

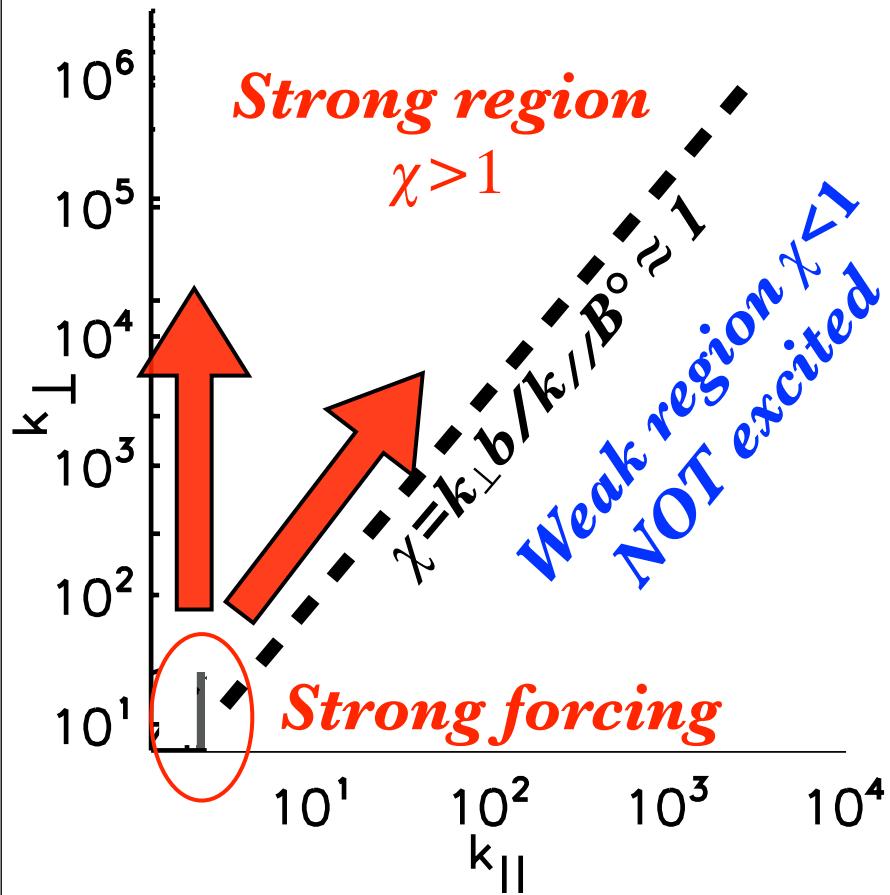
Strong and weak forcing



- In both cases $b_{rms}/B^{\circ} \approx 1/5$

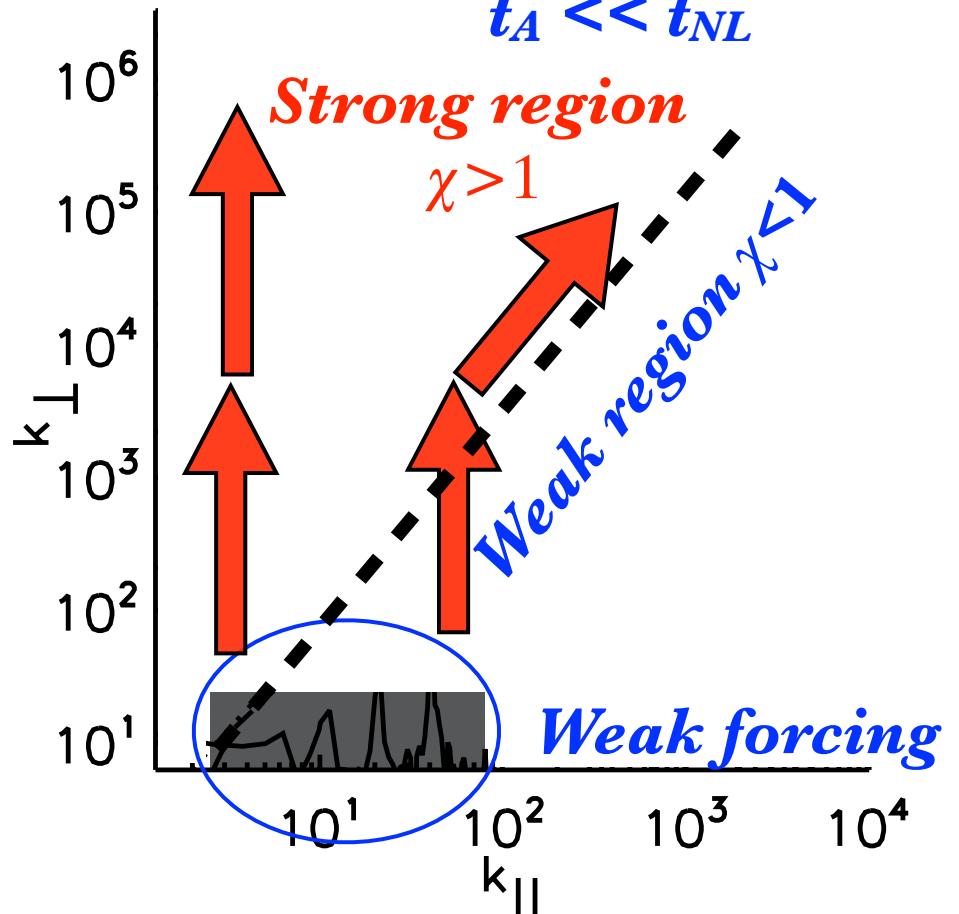
Strong and Weak: Common Wisdom

Enough time : $t_A \approx t_{NL}$



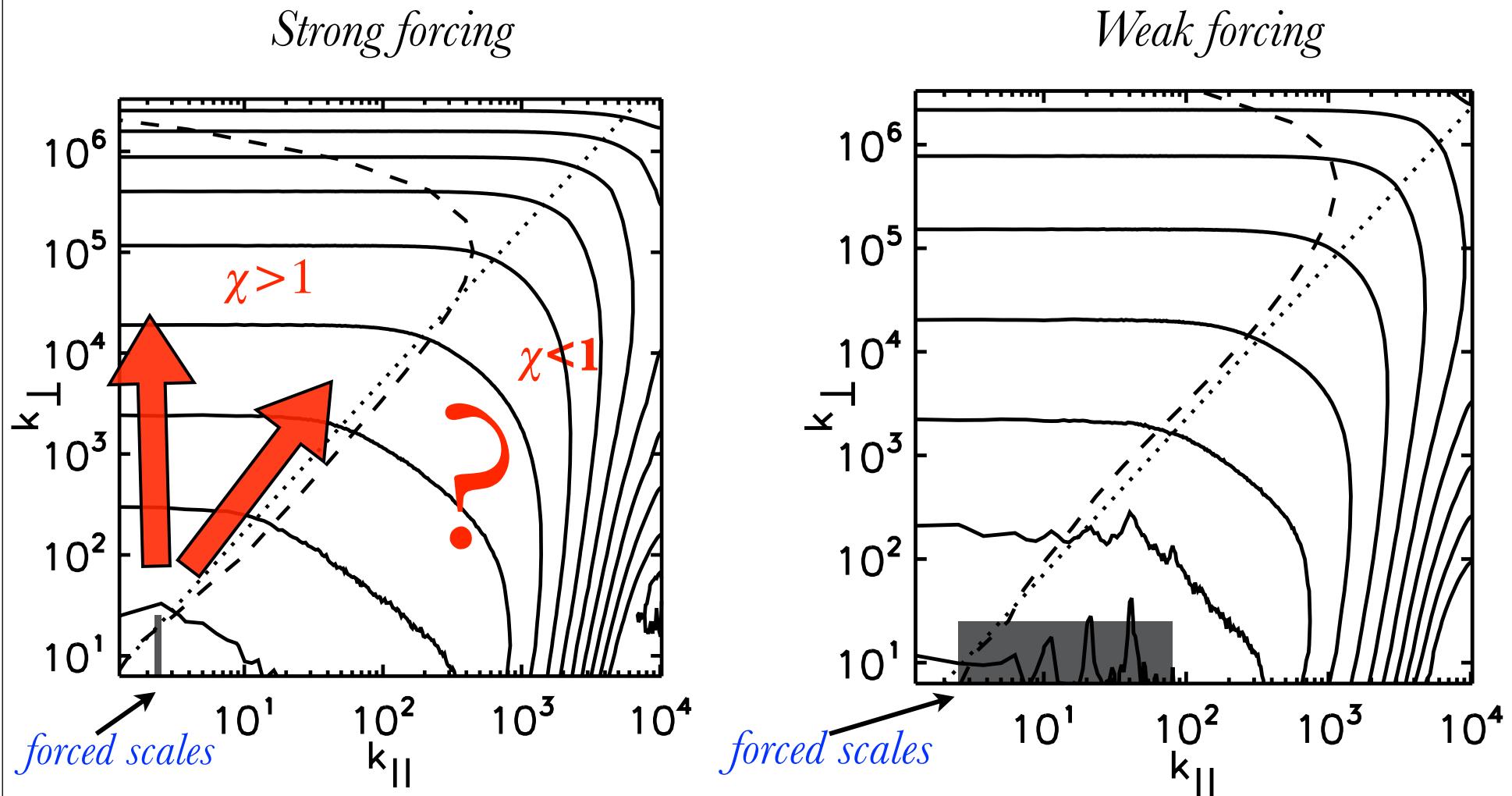
Not enough time

$$t_A \ll t_{NL}$$



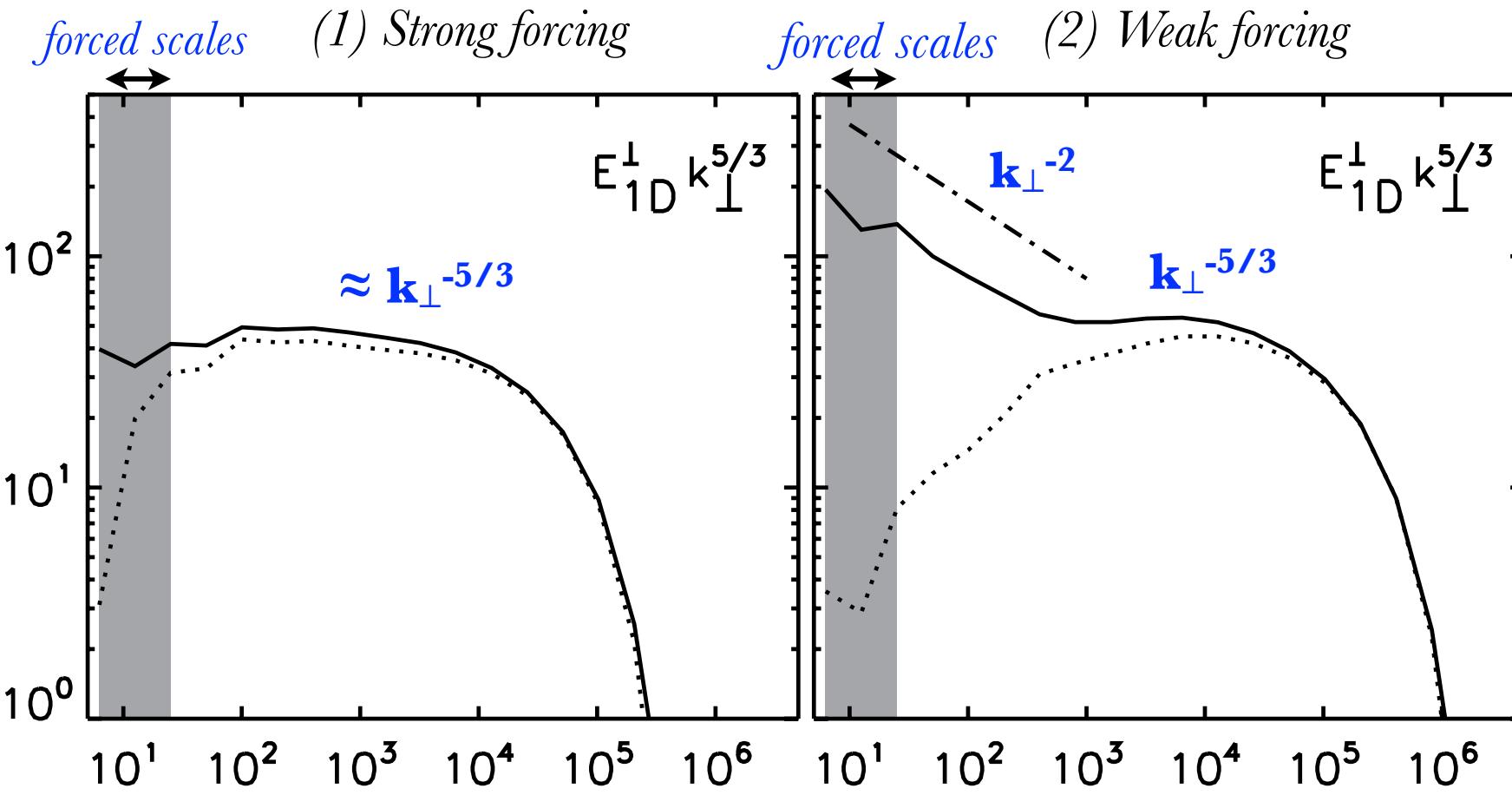
- In both cases $b_{rms}/B^{\circ} \approx 1/5$

Strong and weak forcing: true 3D spectra



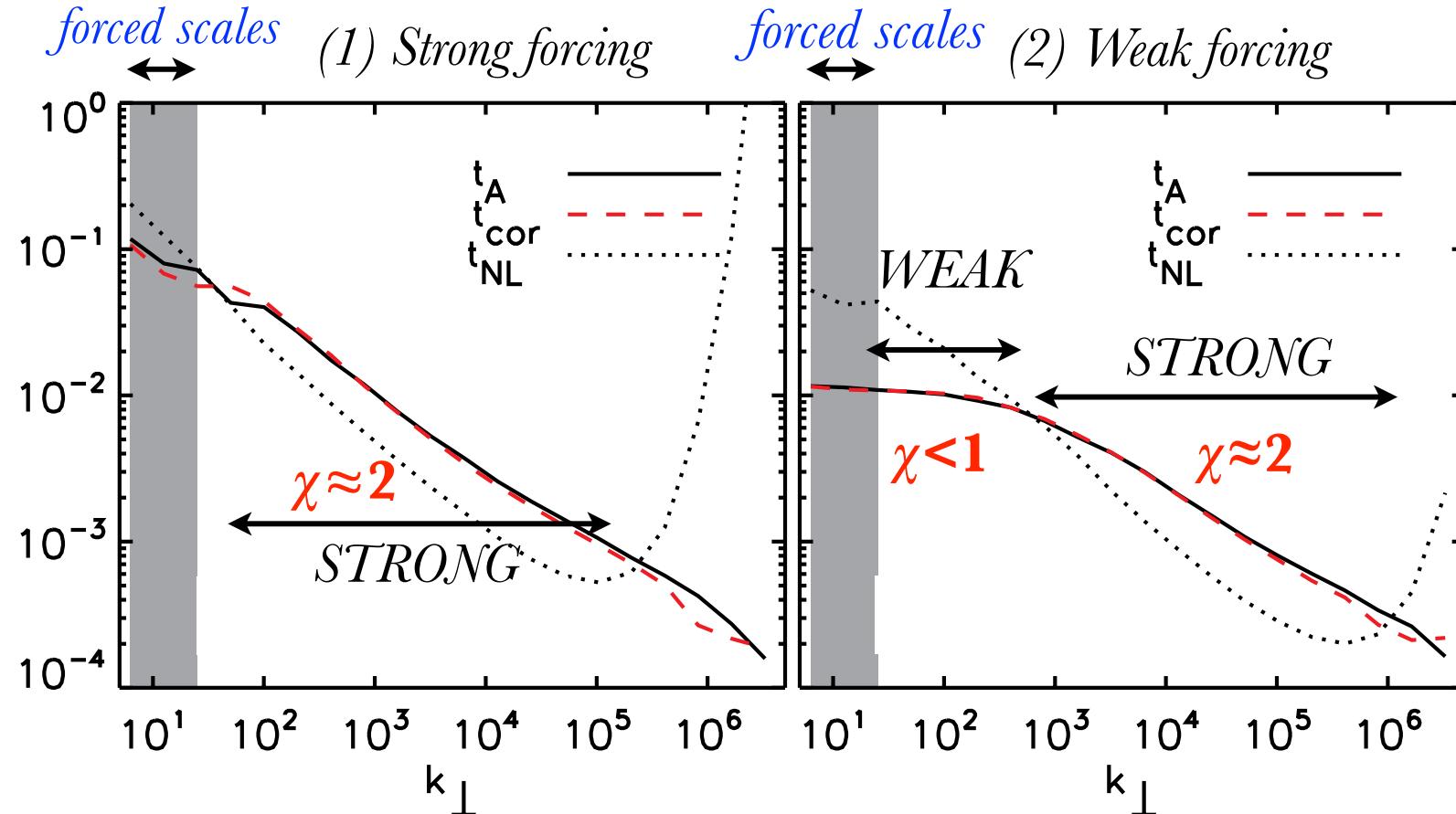
Dashed : $\chi = 1$ isocontour; dotted: theoretical $\chi = 1$ isocontour (neglecting dissipation)

Reduced \perp energy spectrum $E(k_\perp)$

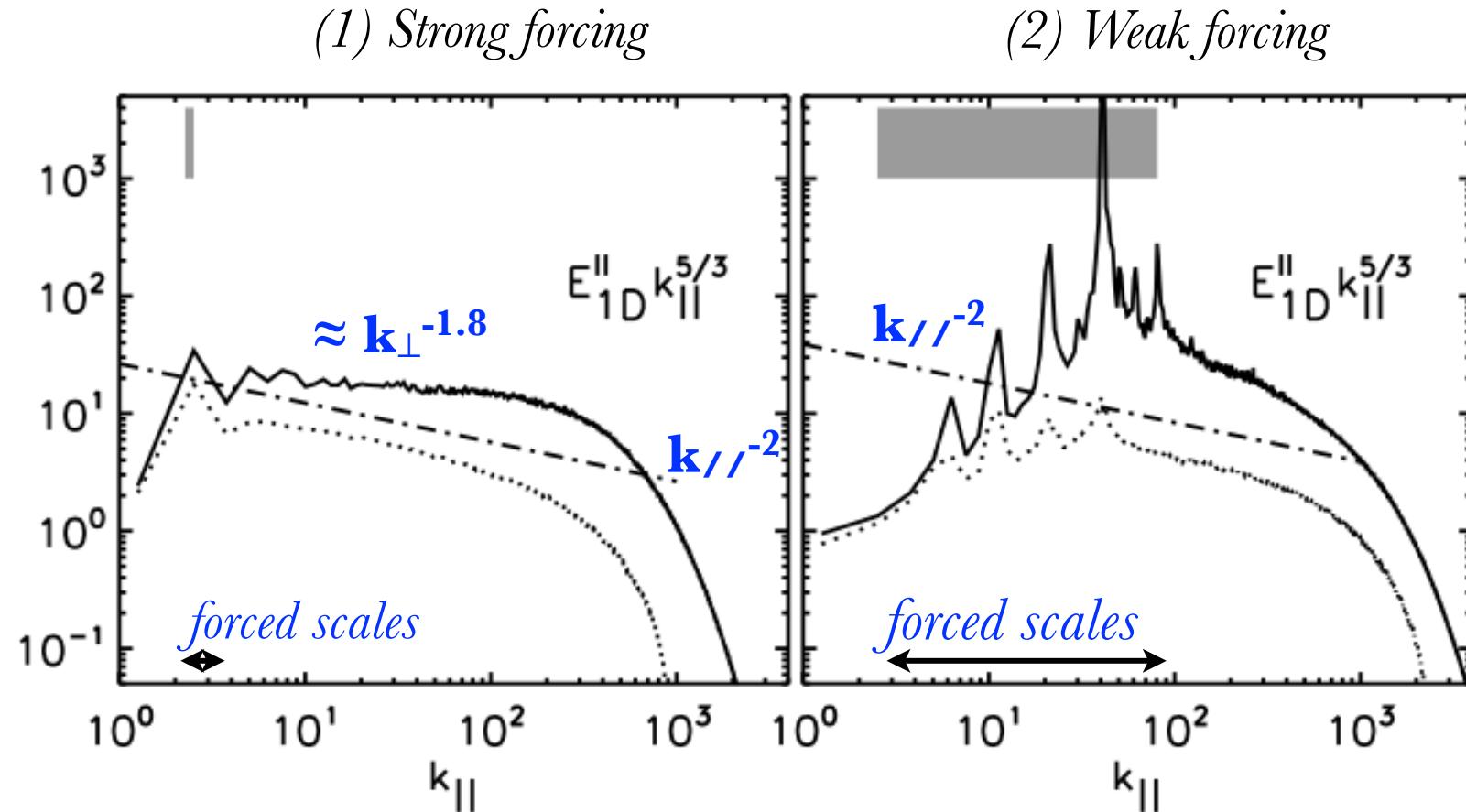


Dotted: reduced spectra obtained from 3D spectra after suppressing excitation in region $\chi < 1/2$
 \Rightarrow this destroys the large scale k_\perp^{-2} in the weak case (right)

Measuring time ratio $\chi=t_A/t_{NL}=1$



Reduced // scaling $E(k_{//})$

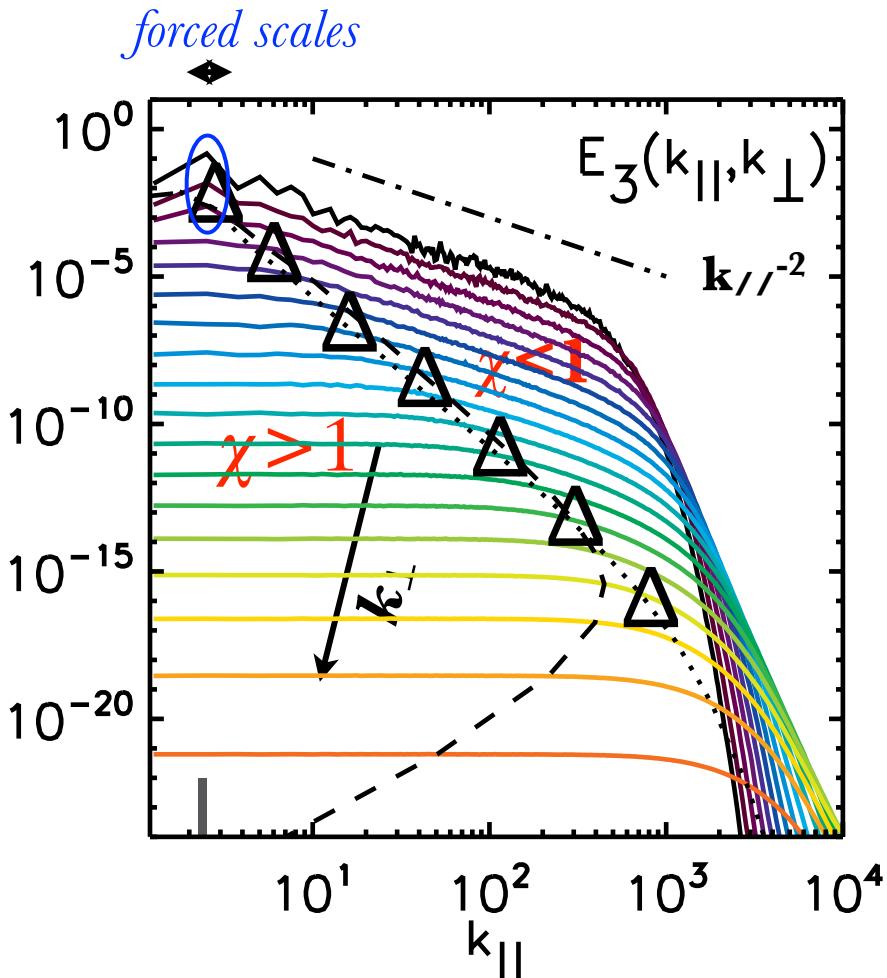


Left: 1D reduced parallel slope is 1.8, NOT -2: this is due to $\chi < 1$ excitation

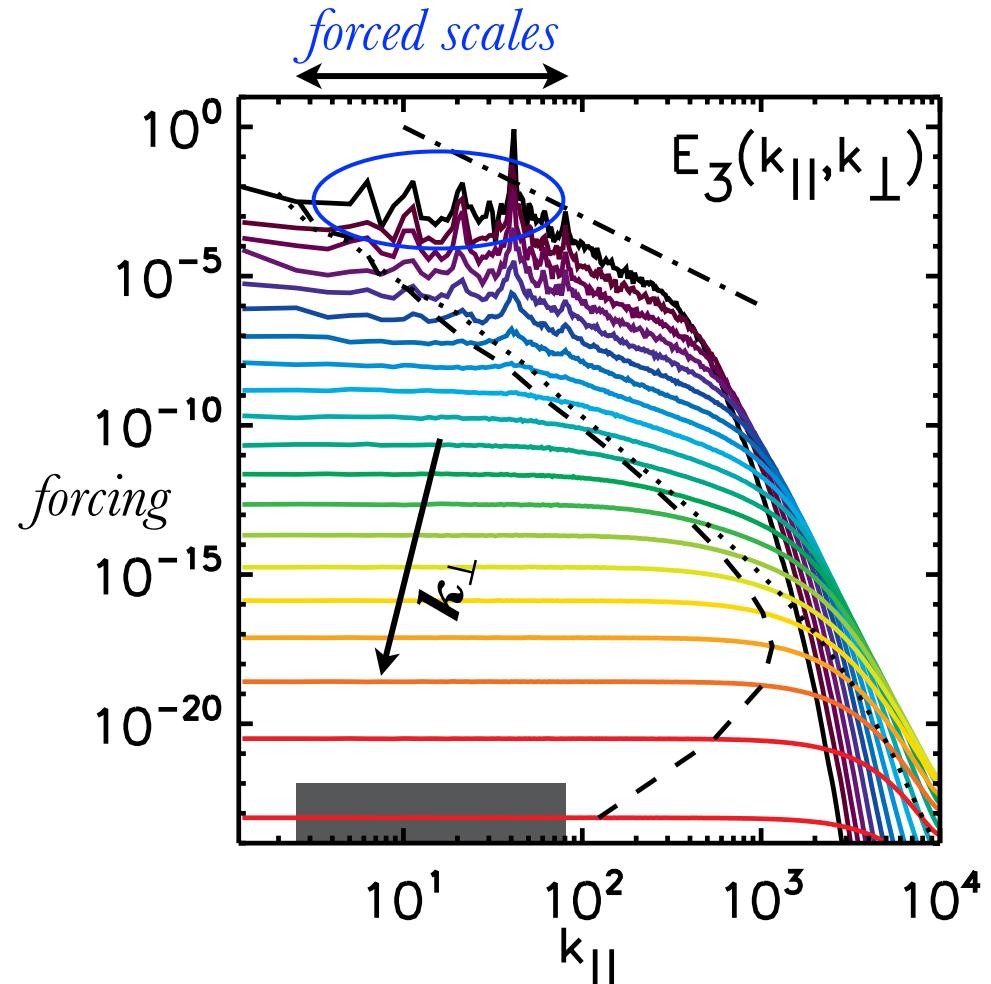
Dotted: reduced spectra obtained from 3D spectra after suppressing excitation in region $\chi < 1/2$

parallel scaling again

(1) Strong forcing



(2) Weak forcing



(b) Energy in $\chi < 1$ region : $\mathbf{E}_3(k_{\perp}, k_{\parallel}) \propto k_{\parallel}^{-2}$

- due to $1/f^2$ spectrum of large \perp eddies
- \neq reduced spectrum $\mathbf{E}_{1D}(k_{\parallel}) \propto k_{\parallel}^{-2}$
- (Δ) marks boundary of "strong" regime $\Delta f = 1/t_{cor}(k_{\perp}) = 1/t_{NL} = kz^{\pm}$

Conclusion

- We obtained here for the first time the double scaling cascade (k_{\perp}^{-2} at large scales then $k_{\perp}^{-5/3}$) of the usual anisotropic phenomenology of RMHD
This result was never obtained in MHD/RMHD perhaps due to too small Reynolds numbers.
- We precisely measured the relevant time scales and found good agreement with both critical balance and weak theory too.
- Interestingly, parallel small scales do get excited in the weak region with a definite scaling. This scaling corresponds to a f^{-2} spectrum providing a tail of harmonics to the turnover time in the strong regime.
The scaling is different when using weak forcing.
- Using the same model and adding large-scale stratification and the associated imbalance between the two Alfvén species leads to different scaling, which might be the source of the large scale $1/f$ spectrum observed in the solar wind (see poster by Grappin Verdini Velli).